

EXAMPLE : COMPLEX ϕ

$$\mathcal{L} = \mathcal{L}(\phi_1) + \mathcal{L}(\phi_2)$$

CHANGE FIELD DEFINITION

$$\phi = \phi_1 + i\phi_2$$

$$\mathcal{L} = \partial_\mu \phi \partial^\mu \phi^* - m^2 \phi \phi^*$$

$$\pi = \dot{\phi}^* \quad \pi^* = \dot{\phi}$$

$$\phi = \int \frac{d^3p}{(2\pi)^3} \left(a_p^\dagger e^{i\vec{p}\cdot\vec{x}} + b_p^\dagger e^{-i\vec{p}\cdot\vec{x}} \right)$$

$$\phi^* = \int \left(a_p^\dagger e^{-i\vec{p}\cdot\vec{x}} + b_p^\dagger e^{i\vec{p}\cdot\vec{x}} \right)$$

$$\mathcal{H} = \int d^3x \left(\pi \dot{\phi} + \pi^* \dot{\phi}^* - \mathcal{L} \right)$$

$$= \int d^3x \left(\pi \pi^* + \vec{\nabla} \phi \cdot \vec{\nabla} \phi^* + m^2 \phi \phi^* \right)$$

$$\dots = \int \frac{d^3p}{(2\pi)^3} E_p \left(a_p^\dagger a_p + b_p^\dagger b_p \right) + C$$

$$j^\mu = i[\partial^\mu \phi^* \phi - \phi^* \partial^\mu \phi]$$

$$Q = \int d^3x j^0 = \int \frac{d^3p}{(2\pi)^3} \left(a_p^\dagger a_p - b_p^\dagger b_p \right)$$

$Q > 0$ $Q < 0$

4. FREE FIELDS

- COMPLEX SCALAR FIELD
- DIRAC FIELD
- GAUGE FIELD

NORMAL ORDERING

$$H = :H: + C$$

::
 { CREATION TO THE LEFT
 DESTRUCTION TO THE RIGHT

HEISENBERG PICTURE:

ψ FIXED
 $i\dot{\phi} = -[H, \phi]$ } C DOES NOT MATTER!

SCHRÖDINGER PICTURE

$$\psi_S = e^{-iHt} \psi_H \rightarrow \psi_S \text{ DOES NOT EXIST!}$$

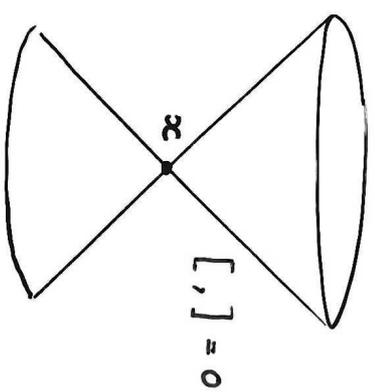
$$O_S = e^{-iHt} O_H e^{iHt}$$

\Rightarrow IN ORDER TO WORK IN EITHER PICTURE, NORMAL-ORDER ALL OPERATORS

CAUSALITY

$$[\phi(x), \phi^*(y)] = \int \frac{d^3p}{2(2\pi)^3 p_0} (e^{-ip \cdot (x-y)} - e^{ip \cdot (y-x)})$$

$$= D(x-y) - D(y-x)$$



$(x-y)^2 < 0$: in the frame $x^0-y^0=0$, rotate $\vec{x}-\vec{y}$ into $\vec{t}-\vec{x}$

$$\Rightarrow [\phi(x), \phi^*(y)] = 0 \text{ if } x^2 + y^2 - 2xy < 0$$

$$\underbrace{\langle 0 | \phi(x) \phi(x)^* | 0 \rangle}_{\text{particle}} - \underbrace{\langle 0 | \phi(x)^* \phi(x) | 0 \rangle}_{\text{antiparticle}}$$

PROBLEM 1 : DIRAC FIELD

$$\mathcal{L} = \bar{\psi} (i\gamma \cdot \partial - m) \psi$$

$$\pi = i\psi^\dagger$$

$$H = \int d^3x \psi^\dagger (-i\vec{\nabla} \cdot \vec{\alpha} - m\beta) \psi$$

$$[\psi_a(\vec{x}), \psi_b^\dagger(\vec{y})] = \delta^{(3)}(\vec{x} - \vec{y}) \delta_{ab}$$

$$\psi(x) = \int \frac{d^3p}{(2\pi)^3 \sqrt{2E_p}} \sum_s (a_p^s u_s e^{i\vec{p} \cdot \vec{x}} + b_p^s v_s e^{-i\vec{p} \cdot \vec{x}})$$

→ WRONG ALGEBRA:

$$[a^r, a^{s\dagger}] = [b^{r\dagger}, b^s] = (2\pi)^3 \delta^{(3)} \delta_{rs}$$

$$\rightarrow b \leftrightarrow b^\dagger$$

H NOT POSITIVE DEFINITE

$$H = \int \frac{d^3p}{(2\pi)^3} \sum E_p (a^\dagger a - b^\dagger b)$$

$$[\psi(x), \bar{\psi}(x)] = \langle 0 | \psi(x) \bar{\psi}(x) | 0 \rangle$$

$$\text{AS } \psi(x) | 0 \rangle = 0$$

→ NEED $E < 0$ TO HAVE CAUSALITY!

SOLUTION

CHANGE THE ALGEBRAS:

$$[,] \rightarrow \{, \}$$

$$\{\psi_a(x), \psi_b^\dagger(y)\}_{ET} = \delta^{(3)}(\vec{x} - \vec{y}) \delta_{ab}$$

$$\{\psi, \psi\}_{ET} = \{\psi^\dagger, \psi^\dagger\}_{ET} = 0$$

$$\psi = \int \frac{d^3p}{(2\pi)^3 \sqrt{2E_p}} \sum a_p^s u^s(\vec{p}) e^{-i\vec{p} \cdot \vec{x}} + b_p^s v^s(\vec{p}) e^{i\vec{p} \cdot \vec{x}}$$

$$\bar{\psi} = \int \frac{d^3p}{(2\pi)^3 \sqrt{2E_p}} \sum b_p^s \bar{v}^s e^{-i\vec{p} \cdot \vec{x}} + a_p^s \bar{u}^s(\vec{p}) e^{i\vec{p} \cdot \vec{x}}$$

$$\Rightarrow \left\{ \begin{array}{l} \{a_p^r, a_q^s\} = \{b_p^r, b_q^s\} = (2\pi)^3 \delta^{(3)}(\vec{p} - \vec{q}) \delta_{rs} \\ \{a^r, a^s\} = \{b^r, b^s\} = 0 \\ \{a_s^\dagger, a_t^\dagger\} = \{b_s^\dagger, b_t^\dagger\} = 0 \end{array} \right.$$

$$H = \int \frac{d^3p}{(2\pi)^3} \sum E_p (a_p^\dagger a_p + b_p^\dagger b_p)$$

CONSEQUENCES

$$\{a_i^+, a_s^+\} = 0 \Rightarrow 1 \text{ or } 0$$

PARTICLE / STATE

$$\{a^+, a\} = \{a, a^+\}$$

SYMMETRY $a \leftrightarrow a^+$
hole particle

OBSERVABLES MUST CONTAIN BIENERS

$$\{\psi(x), \psi(y)\} \text{ VANISHES IF } (x-y)^2 < 0$$

$$\text{BUT WE WANT } [O(x), O(y)] = 0$$

$$[ab, cd] = a\{b\{c\}d\} - a\{c\{b\}d\} + \{a\{c\}d\}b - c\{a\}d\}b$$

SPIN-STATISTICS THEOREM:

$E > 0$, CAUSALITY

$$\Rightarrow \begin{cases} \text{spin } \frac{1}{2} + n & \leftrightarrow \text{fermions} \\ \text{spin } n & \leftrightarrow \text{bosons} \end{cases}$$

PROBLEM 2: GAUGE BOSONS

$$\mathcal{L} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} = -\frac{1}{4} (\partial^\mu A^\nu - \partial^\nu A^\mu)(\partial_\mu A_\nu - \partial_\nu A_\mu)$$

$$\rightarrow \pi^0 = 0$$

$$[A^0, \pi^0] \sim \delta^{(3)} \quad ??$$

\rightarrow • THROW A^0 AWAY AND KEEP ONLY \vec{A}

e.g. $A^0 = 0$ AS GAUGE

• CHANGE \mathcal{L} + RESTRICT π Stueckelberg .

$$\mathcal{L} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} - \frac{1}{2} (\partial_\mu A^\mu)^2$$

$$\Rightarrow \square A^\mu - (1-\lambda) \partial^\mu (\partial \cdot A) = 0$$

$$\left\{ \begin{array}{l} \square A^\mu - (1-\lambda) \partial^\mu (\partial \cdot A) = 0 \\ \pi^\mu = F^{\mu 0} - \lambda g^{\mu 0} (\partial \cdot A) \end{array} \right.$$

$$[A_\rho, A_\nu] = [\pi_\rho, \pi_\nu] = 0$$

$$\lambda=1 \quad \pi_0 = -\partial \cdot A = -\dot{A}_0 + \vec{\nabla} \cdot \vec{A}$$

$$\text{(Feynman)} \quad \pi_i = \partial_i A_0 - \partial_0 A_i = -\dot{A}_i + \vec{\nabla} A^0$$

$$\Rightarrow [\dot{A}_\rho, \dot{A}_\nu] = 0$$

$$[\dot{A}_\rho(x), A_\nu^{\text{ret}}(y)] = i g_{\rho\nu} \delta^{(3)}(\vec{x}-\vec{y})$$

≈ 4 SCALAR FIELDS

$$A_{\mu}(x) = \int \frac{d^3k}{(2\pi)^3 2k} \sum_{\lambda} [a^{\lambda}(k) \epsilon_{\mu}^{\lambda} e^{-ik \cdot x} + a^{\lambda\dagger} \epsilon_{\mu}^{\lambda\dagger} e^{ik \cdot x}]$$

$\partial A^{\mu} = 0 \Rightarrow$ 4 POLARISATIONS

$$\epsilon^{(0)} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \quad \vec{k} \parallel z$$

$$\epsilon^{(1)} = \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} \quad \epsilon^{(2)} = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}$$

$$\Rightarrow [a^{\lambda}, a^{\lambda'\dagger}] = -g^{\lambda\lambda'} \delta^{(3)}(\vec{k} - \vec{k}') (2\pi)^3$$

SHOW THIS

PROBLEM :

$$|1\rangle = a^{0\dagger} |0\rangle$$

$$\langle 1|1\rangle = \langle 0|a^0 a^{0\dagger}|0\rangle$$

$= -\langle 0|0\rangle$ INDEFINITE METRIC!

\Rightarrow SELECT $|1\rangle$ SUCH AS

$$\partial_{\mu} A_{\mu}^{\dagger} |1\rangle = 0$$

$$|1\rangle = |\psi_T\rangle |\phi\rangle \quad \text{AS } i\partial_{\mu} A^{\mu} \sim \sum_{\lambda=0,3} a^{\lambda} \epsilon^{\lambda} \cdot k$$

↑ Transverse
← time + longitudinal

\Rightarrow WE MUST HAVE

$$[a^0 - a^3] |\phi\rangle = 0$$

$$a_3^{\dagger} |0\rangle = |\phi\rangle$$

$$\langle \phi | \phi \rangle = \langle 0 | a_3 a_3^{\dagger} | 0 \rangle$$

$$= \langle 0 | a_0 a_3^{\dagger} | 0 \rangle = 0$$

\Rightarrow STATES OF ZERO NORM

THEY DO NOT ENTER OBSERVABLES:

$$H = \int d^3x \pi^{\mu} \dot{A}_{\mu} - \mathcal{L}$$

$$= \frac{1}{2} \int d^3x \sum_i \dot{A}_i^2 + (\vec{\nabla} \cdot \vec{A})^2 - A_0^2 - (A_3)^2$$

$$= \int \frac{d^3k}{(2\pi)^3} \epsilon_k \sum_{\lambda=1}^3 (a^{\lambda\dagger} a^{\lambda}) - a^{0\dagger} a^0$$

SHOW THIS

$\langle \psi | H | \psi \rangle$ INDEPENDENT OF a_0, a_3

Δ TRUE FOR ON-SHELL PHOTONS

OFF-SHELL: $\partial \cdot A \rightarrow k \cdot \epsilon$

$$g_{\mu\nu} = \sum_{\lambda} \frac{\epsilon_{\mu}^{(\lambda)} \epsilon_{\nu}^{(\lambda)}}{k \cdot \epsilon^{(\lambda)}}$$

$\rightarrow \epsilon_3$ "SURVIVES"

- CAUSALITY
- LORENTZ INVARIANT
- RENORMALIZABLE
- EXTRA SYMMETRIES: GAUGE, ETC.

$$\mathcal{L} = \mathcal{L}_0 + \mathcal{L}_I$$

HOW DO THESE INTERACT

GAUGE FIELDS	FERMIONS	SCALAR BOSONS
$-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{z}{2} (a \cdot A)^2$ $\partial_{\mu} F^{\mu\nu} \text{ broken}$ $\pi_0 = a \cdot A$ $\pi_i = a_i A_0 - a_0 A_i$ $\left. \begin{matrix} a_+^+ \\ a_0, a_3^+ \\ a_+^- \end{matrix} \right\} + \rangle \rightarrow 0$	$\bar{\psi} (i \partial \cdot \gamma - m) \psi$ $\bar{\psi} \gamma_{\mu} \psi$ $\pi = i \psi^+ \psi$ $\pi^+ = -i \psi \psi^+$ $\psi \sim a e^{-i p \cdot x} + b^+ e^{i p \cdot x}$ $\bar{\psi} \sim a' a^+ + b' b^+ \delta_{rs}$	$\mathcal{L} = \partial_{\mu} \phi \partial^{\mu} \phi^* - m^2 \phi \phi^*$ $\partial_{\mu} = i(\phi^* \partial_{\mu} \phi - \phi \partial_{\mu} \phi^*)$ $\pi = \dot{\phi}^*$ $\pi^* = \dot{\phi}$ $\phi \sim a e^{-i p \cdot x} + b^+ e^{i p \cdot x}$ $[a, a^+] = (2\pi)^3 \delta^{(3)} = [b, b^+]$