

4. FREE FIELDS

- COMPLEX SCALAR FIELD
- DIRAC FIELD
- GAUGE FIELD

EXAMPLE : COMPLEX ϕ

$$\mathcal{L} = \mathcal{L}(\phi_1) + \mathcal{L}(\phi_2)$$

CHANGE FIELD DEFINITION

$$\phi = \phi_1 + i\phi_2$$

$$\mathcal{L} = \partial_\mu \phi \partial^\mu \phi^* - m^2 \phi \phi^*$$

$$\pi = \dot{\phi}^* \quad \pi^* = \dot{\phi}$$

$$\phi = \int \frac{d^3p}{\sqrt{2\pi} (2\pi)^3} \left(a_{\vec{p}} e^{i\vec{p}\cdot\vec{x}} + b_{\vec{p}}^\dagger e^{-i\vec{p}\cdot\vec{x}} \right)$$

$$\phi^* = \int \frac{d^3p}{\sqrt{2\pi} (2\pi)^3} \left(a_{\vec{p}}^\dagger e^{-i\vec{p}\cdot\vec{x}} + b_{\vec{p}} e^{i\vec{p}\cdot\vec{x}} \right)$$

$$\begin{aligned} \mathcal{H} &= \int d^3x \left(\pi \dot{\phi} + \pi^* \dot{\phi}^* - \mathcal{L} \right) \\ &= \int d^3x \left(\pi \pi^* + \vec{\nabla} \phi \cdot \vec{\nabla} \phi^* + m^2 \phi \phi^* \right) \\ &\dots \\ &= \int \frac{d^3p}{(2\pi)^3} E_p \left(a_{\vec{p}}^\dagger a_{\vec{p}} + b_{\vec{p}}^\dagger b_{\vec{p}} \right) + C \end{aligned}$$

$$j^\mu = i \left[(\partial^\mu \phi^*) \phi - \phi^* \partial^\mu \phi \right]$$

$$Q = \int d^3x j^0 = \int \frac{d^3p}{(2\pi)^3 (2\pi)} \left(\underset{\uparrow}{a_{\vec{p}}^\dagger} a_{\vec{p}} - b_{\vec{p}}^\dagger \underset{\uparrow}{b_{\vec{p}}} \right)$$

$Q > 0 \qquad \qquad \qquad Q < 0$

NORMAL ORDERING

$$H = :H: + C$$

$::$
 { CREATION TO THE LEFT
 DESTRUCTION TO THE RIGHT

HEISENBERG PICTURE:

$$\left. \begin{array}{l} \psi \text{ FIXED} \\ i\dot{O} = -[H, O] \end{array} \right\} C \text{ DOES NOT MATTER!}$$

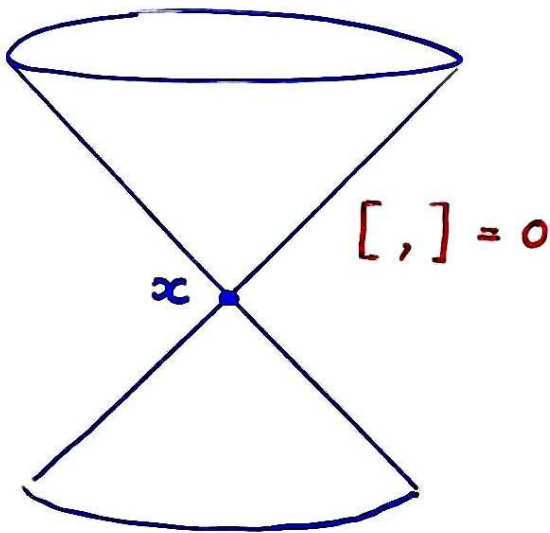
SCHRÖDINGER PICTURE

$$\begin{array}{l} \psi_S = e^{-Ht} \psi_H \\ O_S = e^{-iHt} O_H e^{iHt} \end{array} \rightarrow \psi_S \text{ DOES NOT EXIST!}$$

\Rightarrow IN ORDER TO WORK IN EITHER PICTURE, NORMAL-ORDER ALL OPERATORS

CAUSALITY

$$\begin{aligned}
 [\phi(x), \phi^*(y)] &= \int \frac{d^3p}{2(2\pi)^3 p^0} \left(e^{-ip \cdot (x-y)} - e^{ip \cdot (y-x)} \right) \\
 &= D(x-y) - D(y-x)
 \end{aligned}$$



$(x-y)^2 < 0$: in the frame $x^0 - y^0 = 0$,
rotate $\vec{x} - \vec{y}$ into $\vec{y} - \vec{x}$

$$\Rightarrow [\phi(x), \phi^*(y)] = 0 \quad \text{if} \quad x^2 + y^2 - 2xy < 0$$

$$\underbrace{\langle 0 | \phi(x) \phi(x)^* | 0 \rangle}_{\text{particle}} - \underbrace{\langle 0 | \phi(x)^* \phi(x) | 0 \rangle}_{\text{antiparticle}}$$

PROBLEM 1 : DIRAC FIELD

$$\mathcal{L} = \bar{\psi} (i \gamma \cdot \partial - m) \psi$$

$$\pi = i \psi^\dagger$$

$$H = \int d^3x \psi^\dagger (-i \vec{\sigma} \cdot \vec{\alpha} - m \beta) \psi$$

$$[\psi_a(\vec{x}), \psi_b^\dagger(\vec{y})] = \delta^{(3)}(\vec{x} - \vec{y}) \delta_{ab}$$

$$\psi(x) = \int \frac{d^3p}{(2\pi)^3 \sqrt{2E_p}} \sum_s (a_{\vec{p}}^s u_s e^{i\vec{p} \cdot \vec{x}} + b_{\vec{p}}^{s\dagger} v_s e^{-i\vec{p} \cdot \vec{x}})$$

→ WRONG ALGEBRA:

$$[a^r, a^{s\dagger}] = [b^{tr}, b^s] = (2\pi)^3 \delta^{(3)} \delta_{rs}$$

→ $b \leftrightarrow b^\dagger$

H NOT POSITIVE DEFINITE

$$H = \int \frac{d^3p}{(2\pi)^3} \sum E_p (a^\dagger a - \underset{\uparrow}{b^\dagger b})$$

$$[\psi(x), \bar{\psi}(x)] = \langle 0 | \psi(x) \bar{\psi}(y) | 0 \rangle$$

$$\text{AS } \psi(x) | 0 \rangle = 0$$

→ NEED $E < 0$ TO HAVE CAUSALITY!

SOLUTION

CHANGE THE ALGEBRAS:

$$[,] \rightarrow \{ , \}$$

$$\{ \psi_a(x), \psi_b^\dagger(y) \}_{ET} = \delta^{(3)}(\vec{x} - \vec{y}) \delta^{ab}$$

$$\{ \psi, \psi \}_{ET} = \{ \psi^\dagger, \psi^\dagger \}_{ET} = 0$$

$$\psi = \int \frac{d^3p}{(2\pi)^3 \sqrt{2E_p}} \sum a_p^s u^s(\vec{p}) e^{-ip \cdot x} + b_p^{s\dagger} v^s(\vec{p}) e^{ip \cdot x}$$

$$\bar{\psi} = \int \frac{d^3p}{(2\pi)^3 \sqrt{2E_p}} \sum b_p^s \bar{v}^s e^{-ip \cdot x} + a_p^{s\dagger} \bar{u}^s(\vec{p}) e^{ip \cdot x}$$

$$\Rightarrow \begin{cases} \{ a_p^r, a_q^{s\dagger} \} = \{ b_p^s, b_q^{s\dagger} \} = (2\pi)^3 \delta^{(3)}(\vec{p} - \vec{q}) \delta_{rs} \\ \{ a^r, a^s \} = \{ b^r, b^s \} = 0 \\ \{ a_s^\dagger, a_r^\dagger \} = \{ b_s^\dagger, b_r^\dagger \} = 0 \end{cases}$$

$$H = \int \frac{d^3p}{(2\pi)^3} \sum_p E_p (a_p^{s\dagger} a_p^s + b_p^{s\dagger} b_p^s)$$

CONSEQUENCES

$$\{a_r^+, a_s^+\} = 0 \quad \Rightarrow \quad 1 \propto 0$$

PARTICLE / STATE

$$\{a^+, a\} = \{a, a^+\}$$

SYMMETRY $a \rightarrow a^+$
hole particle

OBSERVABLES MUST CONTAIN BILINEARS

$$\{\psi(x), \psi(y)\} \text{ VANISHES IF } (x-y)^2 < 0$$

$$\text{BUT WE WANT } [O(x), O(y)] = 0$$

$$[ab, cd] = a\{bc\}d - ac\{bd\} + \{ac\}db - c\{ad\}b$$

SPIN-STATISTICS THEOREM :

$E > 0$, CAUSALITY

$$\Rightarrow \begin{cases} \text{spin } \frac{1}{2} + n \\ \text{spin } n \end{cases} \begin{matrix} \leftrightarrow \text{ fermions} \\ \leftrightarrow \text{ bosons} \end{matrix}$$

PROBLEM 2 : GAUGE BOSONS

$$\mathcal{L} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} = -\frac{1}{4} (\partial^\mu A^\nu - \partial^\nu A^\mu) (\partial_\mu A_\nu - \partial_\nu A_\mu)$$

$$\rightarrow \pi^0 = 0$$

$$[A^0, \pi^0] \sim \delta^{(3)} \quad ??$$

\rightarrow • THROW A^0 AWAY AND KEEP ONLY \vec{A}

e.g. $A^0 = 0$ AS GAUGE

• CHANGE \mathcal{L} + RESTRICT Hilbert sp.

$$\mathcal{L} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu} - \frac{\lambda}{2} (\partial_\mu A^\mu)^2$$

$$\Rightarrow \begin{cases} \square A^\mu - (1-\lambda) \partial^\mu (\partial \cdot A) = 0 \\ \pi^\mu = F^{\mu 0} - \lambda g^{\mu 0} (\partial \cdot A) \end{cases}$$

$$[A_\rho, A_\nu] = [\pi_\rho, \pi_\nu] = 0$$

$$\lambda=1 \quad \begin{aligned} \pi_0 &= -\partial \cdot A = -\dot{A}_0 + \vec{\nabla} \cdot \vec{A} \\ \pi_i &= \partial_i A_0 - \partial_0 A_i = -\dot{A}_i + \vec{\nabla} A_0 \end{aligned}$$

(Feynman)

$$\Rightarrow [\dot{A}_\rho, \dot{A}_\nu] = 0$$

$$[\dot{A}_\rho(x), A_\nu^{\text{ret}}(\vec{y}, t)] = i g_{\rho\nu} \delta^{(3)}(\vec{x} - \vec{y})$$

\approx 4 SCALAR FIELDS

$$A_\mu(x) = \int \frac{d^3k}{(2\pi)^3 \sqrt{2k_0}} \sum_\lambda [a^\lambda(k) \epsilon_\mu^\lambda e^{-i\vec{k}\cdot\vec{x}} + a^{\lambda\dagger} \epsilon_\mu^{\lambda*} e^{i\vec{k}\cdot\vec{x}}]$$

$\square A^\mu = 0 \Rightarrow$ 4 POLARISATIONS

$$\epsilon^{(0)} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix} \quad \epsilon^{(3)} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \end{pmatrix} \quad \vec{k} \parallel z$$

$$\epsilon^{(1)} = \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} \quad \epsilon^{(2)} = \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}$$

$$\Rightarrow [a^\lambda, a^{\lambda'\dagger}] = -g^{\lambda\lambda'} \delta^{(3)}(\vec{k}-\vec{k}') (2\pi)^3$$

SHOW THIS

PROBLEM :

$$|1\rangle = a^{0\dagger} |0\rangle$$

$$\langle 1|1\rangle = \langle 0|a^0 a^{0\dagger}|0\rangle$$

$$= -\langle 0|0\rangle$$

INDEFINITE METRIC!

\Rightarrow SELECT $|\psi\rangle$ SUCH AS

$$\partial^\mu A_\mu^\dagger |\psi\rangle = 0$$

$$|\psi\rangle = |\psi_T\rangle |\phi\rangle$$

↑
transverse

↑
time + longitudinal

$$\text{AS } i\partial \cdot A^\dagger \sim \sum_{\lambda=03} a^\lambda \epsilon^\lambda \cdot \vec{k}$$

⇒ WE MUST HAVE

$$[a^0 - a^3] |\phi\rangle = 0$$

$$a_3^\dagger |0\rangle = |\phi\rangle$$

$$\begin{aligned} \langle \phi | \phi \rangle &= \langle 0 | a_3 a_3^\dagger | 0 \rangle \\ &= \langle 0 | a_0 a_3^\dagger | 0 \rangle = 0 \end{aligned}$$

⇒ STATES OF ZERO NORM

THEY DO NOT ENTER OBSERVABLES:

$$\begin{aligned} H &= \int d^3x \pi^k \dot{A}_k - \mathcal{L} \\ &= \frac{1}{2} \int d^3x \sum_{i=1}^3 \dot{A}_i^2 + (\vec{\nabla} A_i)^2 - A_0^2 - (\vec{\nabla} A_0)^2 \\ &= \int \frac{d^3k}{(2\pi)^3} E_k \sum_{\lambda=1}^3 (a^{\lambda\dagger} a^\lambda) - a^{0\dagger} a^0 \end{aligned}$$

SHOW THIS

$\langle \psi | H | \psi \rangle$ INDEPENDENT OF a_0, a_3

△ TRUE FOR ON-SHELL PHOTONS

OFF-SHELL: $\partial \cdot A \rightarrow k \cdot \epsilon$

$$g_{\mu\nu} = \sum_{\lambda} \frac{\epsilon_{\mu}^{(\lambda)} \epsilon_{\nu}^{(\lambda)}}{k \cdot \epsilon^{(\lambda)}}$$

→ ϵ_3 "SURVIVES"

SCALAR BOSONS

$$\mathcal{L} = \partial_\mu \phi \partial^\mu \phi^* - m^2 \phi \phi^*$$

$$d_\mu = i(\phi^* \partial_\mu \phi - \phi \partial_\mu \phi^*)$$

$$\pi = \dot{\phi}^* \quad \pi^* = \dot{\phi}$$

$$\phi \sim a e^{-ip \cdot x} + b^\dagger e^{ip \cdot x}$$

$$[a, a^\dagger] = (2\pi)^3 \delta^{(3)} = [b, b^\dagger]$$

FERMIONS

$$\bar{\psi} (i \not{\partial} \cdot \gamma - m) \psi$$

$$\bar{\psi} \not{\partial}_\mu \psi$$

$$\pi = i \psi^\dagger \quad \pi^\dagger = -i \psi$$

$$\psi \sim a e^{-ip \cdot x} + b^\dagger e^{ip \cdot x}$$

$$\{a, a^\dagger\} = \{b, b^\dagger\} = (2\pi)^3 \delta^{(3)} \delta_{rs}$$

GAUGE FIELDS

$$-\frac{1}{4} F^{\mu\nu} F_{\mu\nu} - \frac{1}{2} (D_\mu A)^2$$

$\partial_\mu F^{\mu\nu}$ broken

$$\pi_0 = D_\mu A$$

$$\pi_i = \partial_i A_0 - \partial_0 A_i$$

$$\left\{ \begin{array}{l} a_0^+, a_j^+ \\ a_+^+ \end{array} \right\} | \psi \rangle \rightarrow 0$$

HOW DO THESE INTERACT

$$\mathcal{L} = \mathcal{L}_0 + \mathcal{L}_I$$

- CAUSALITY
- LORENTZ INVARIANT
- RENORMALISABLE
- EXTRA SYMMETRIES: GAUGE, ETC.